# Phase Space Analysis for Evolution PDE's and Applications<sup>1</sup> Global existence (in time) for semilinear models

#### Marcelo Rempel Ebert - University of São Paulo

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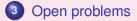
# Summary





2 Semilinear damped wave equation

- Non-existence result
- Global existence results



# A general problem

Let us consider the semilinear Cauchy problem

$$L(\partial_t,\partial_x,t,x)u=f(u), \ u(0,x)=u_0(x), \ u_t(0,x)=u_1(x),$$

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<sup>&</sup>lt;sup>2</sup>M. R. Ebert, M. Reissig, Methods for Partial Differential Equations, Qualitative properties of solutions - Phase space analysis - Semi-linear models, Birkhäuser, Cham (2018).

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Depending on the regularity of initial data,  $p = p_c$  belongs to the supercritical or subcritical case.

# Semilinear heat model

Let us consider the semilinear heat model with power nonlinearity

$$u_t - \triangle u = |u|^{p-1}u, \ u(0, x) = u_0(x).$$

Fujita <sup>3</sup> shown that  $p_{Fuj}(n) = 1 + \frac{2}{n}$  is the threshold between global existence of small data solutions for exponents larger and blow up behavior of solutions for exponents smaller or equal to the Fujita exponent.

<sup>3</sup>H. Fujita, On the blowing up of solutions of the Cauchy problem for  $u_t = \triangle u + u^{1+\alpha}$ . J. Fac. Sci. Univ. Tokyo Sect. I 13 (1966), 109–124.

# Semilinear wave equation

Much has been devoted to the Cauchy problem

$$\begin{cases} u_{tt} + (-\triangle)^{\sigma} u = |u|^{\rho}, & t \ge 0, \ x \in \mathbb{R}^n \\ u(0, x) = u_0(x), & u_t(0, x) = u_1(x). \end{cases}$$
(1)

 If σ = 1, it was proved <sup>4</sup> that the critical index p<sub>0</sub>(n) to (1) is the Strauss exponent, i.e., the positive root of

$$(n-1)p^2 - (n+1)p - 2 = 0.$$

 If σ > 1, then the critical exponent to (1) is given by E. & Lourenço (2019)

$$p_c \doteq 1 + \frac{2\sigma}{[n-\sigma]_+}.$$

<sup>4</sup>V. Georgiev, H. Lindblad, and C. Sogge, Weighted Strichartz estimates and global existence for semilinear wave equations. Amer. J. Math. 119 (1997), no. 6, 1291–1319.

# Semilinear exterior damped wave equations: 1/3

A priori estimates for solutions to the associate linear problem <sup>5</sup> may be applied to look for global-in-time small data solutions to

 $u_{tt} - \triangle u + u_t = f(u),$   $(u, u_t)(0, x) = (u_0, u_1)(x),$ 

where  $|f(u)| \approx |u|^p$  for some p > 1.

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where  $|f(u)| \approx |u|^p$  for some p > 1. By using comp. supp. data in  $H^1 \times L^2$ , Todorova & Yordanov (2001) showed that critical exponent is given by Fujita index  $p_c = 1 + 2/n$  for any  $n \ge 1$ :

- Global existence for small data if p > 1 + 2/n (and also  $p \le 1 + 2/(n-2)$  if  $n \ge 3$ ),
- Nonexistence if p < 1 + 2/n (p = 1 + 2/n, in Zhang (2001))</p>

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# Removing compact support: 2/3

• In <sup>6</sup> the authors showed that comp. supp. is not necessary. If data are small in  $H^1 \times L^2$  and in  $L^1$ , critical exponent remains 1 + 2/n in space dimension n = 1, 2. For n = 3, they could proved a Global existence only for 2 .

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- Then Narazaki (2004) introduced estimates on L<sup>p</sup> − L<sup>q</sup> basis, extending to space dimension n ≤ 5 if data are small in: energy space, L<sup>1</sup> and W<sup>1,p</sup> × L<sup>p</sup> when p < 2.</li>

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- Ikehata & Ohta (2002) for n = 1, 2, proved that if one removes the assumption that the initial data are in  $L^1(\mathbb{R}^n)$ and only assume that they are in the energy space, then the critical exponent is modified to  $1 + \frac{2m}{n}$  under additional regularity  $L^m(\mathbb{R}^n)$ , with  $m \in [1, 2]$ .

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<sup>&</sup>lt;sup>7</sup>M. Ikeda, T. Inui, M. Okamoto, Y. Wakasugi,  $L^p - L^q$  estimates for the damped wave equation and the critical exponent for the nonlinear problem with slowly decaying data, Communications on Pure and Applied Analysis, 18 (2019), 1967–2008.

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- E.& Girardi & Reissig (Math. Ann. 2019) considered

$$u_{tt} - \bigtriangleup u + u_t = |u|^{1+\frac{2}{n}} \mu(|u|), \qquad (u, u_t)(0, x) = (u_0, u_1)(x),$$

where  $\mu$  is a modulus of continuity. Under additional regularity  $L^1(\mathbb{R}^n)$  for initial data, the threshold between global (in time) existence of small data solutions and blow-up behavior even of small data solutions is given by

$$\int_0^\epsilon \frac{\mu(s)}{s}\,ds < \infty \ \text{ or } \ \int_0^\epsilon \frac{\mu(s)}{s}\,ds = \infty, \quad \epsilon > 0 \text{ sufficiently small}.$$

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# Non-existence result via Test function method

#### Theorem 1

Let us consider the Cauchy problem for the classical damped wave equation with power nonlinearity

$$u_{tt} - \triangle u + u_t = |u|^p, \ u(0, x) = \varphi(x), \ u_t(0, x) = \psi(x)$$

in  $[0, \infty) \times \mathbb{R}^n$  with  $n \ge 1$  and  $p \in (1, 1 + \frac{2}{n}]$ . Let  $(\varphi, \psi) \in \mathcal{A}_{1,1}$  satisfy the assumption

$$\int_{\mathbb{R}^n} \left( \varphi(x) + \psi(x) \right) \, dx > 0.$$

Then there exists a locally (in time) defined energy solution

$$u \in C([0, T), H^1(\mathbb{R}^n)) \cap C^1([0, T), L^2(\mathbb{R}^n)).$$

This solution can not be continued to the interval  $[0,\infty)$  in time.

# Proof of Theorem 1: 1/4

We first introduce test functions  $\eta = \eta(t)$  and  $\phi = \phi(x)$  having the following properties:

$$\begin{array}{l} \bullet \quad \eta \in C_0^{\infty}[0,\infty), \ 0 \leq \eta(t) \leq 1, \\ \eta(t) = \begin{cases} 1 & \text{for } 0 \leq t \leq \frac{1}{2}, \\ 0 & \text{for } t \geq 1, \end{cases} \\ \begin{array}{l} \bullet \quad \phi(x) \leq \phi(x) \leq 1, \\ \phi(x) = \begin{cases} 1 & \text{for } |x| \leq \frac{1}{2}, \\ 0 & \text{for } |x| \geq 1, \end{cases} \\ \begin{array}{l} \bullet \quad \frac{\eta'(t)^2}{\eta(t)} \leq C \text{ for } \frac{1}{2} < t < 1, \text{ and } \frac{|\nabla \phi(x)|^2}{\phi(x)} \leq C \text{ for } \frac{1}{2} < |x| < 1. \end{cases} \\ \begin{array}{l} \bullet \quad H \in [0,\infty) \text{ be a large parameter. We define the test function} \end{cases}$$

$$\chi_{R}(t, \mathbf{x}) := \eta_{R}(t)\phi_{R}(\mathbf{x}) := \eta\left(\frac{t}{R^{2}}\right)\phi\left(\frac{\mathbf{x}}{R}\right).$$

# Proof of Theorem 1: 2/4

We put

$$Q_R := [0, R^2] \times B_R, \quad B_R := \{x \in \mathbb{R}^n : |x| \le R\}.$$

We suppose that the energy solution u = u(t, x) exists globally in time. We define the functional

$$I_R := \int_{Q_R} |u(t,x)|^p \chi_R(t,x)^q d(x,t) = \int_{Q_R} (u_{tt} - \bigtriangleup u + u_t) \chi_R(t,x)^q d(x,t).$$

Here *q* is the Sobolev conjugate of *p*, that is,  $\frac{1}{p} + \frac{1}{q} = 1$ . After integration by parts we obtain

$$I_{R} = -\int_{B_{R}} (\varphi + \psi) \phi_{R}^{q} dx + \int_{Q_{R}} u \partial_{t}^{2}(\chi_{R}^{q}) d(x, t)$$
$$-\int_{Q_{R}} u \partial_{t}(\chi_{R}^{q}) d(x, t) - \int_{Q_{R}} u \triangle(\chi_{R}^{q}) d(x, t)$$
$$\coloneqq -\int_{B_{R}} (\varphi + \psi) \phi_{R}^{q} dx + J_{1} + J_{2} + J_{3}.$$

# Proof of Theorem 1: 3/4

We shall estimate separately  $J_1$ ,  $J_2$  and  $J_3$ . Here we use the notations

$$\hat{Q}_{R,t} := \left[\frac{R^2}{2}, R^2\right] \times B_R, \quad \hat{Q}_{R,x} := [0, R^2] \times (B_R \setminus B_{\frac{R}{2}}).$$

Application of Hölder's inequality implies

$$\begin{aligned} |J_{3}| & \leq CR^{-2} \Big( \int_{\hat{Q}_{R,x}} |u|^{p} \chi_{R}^{q}(t,x) \, d(x,t) \Big)^{1/p} \Big( \int_{\hat{Q}_{R,x}} 1 \, d(x,t) \Big)^{1/q} \\ & \leq CR^{-2} I_{R,x}^{\frac{1}{p}} \Big( \int_{\hat{Q}_{R,x}} 1 \, d(x,t) \Big)^{1/q} \leq C I_{R,x}^{\frac{1}{p}} R^{\frac{n+2}{q}-2}, \end{aligned}$$

where

$$I_{R,x} := \int_{\hat{Q}_{R,x}} |u|^p \chi^q_R(t,x) \, d(x,t).$$

Since  $1 the last inequality gives <math>|J_3| \le C I_{R,x}^{\frac{1}{p}}$ .

# Proof of Theorem 1: 4/4

By

$$\partial_t(\chi^q_R) = \frac{1}{R^2} q \phi^q_R(x) \eta^{q-1}_R(t) \eta' \Big( \frac{t}{R^2} \Big),$$

we have

$$\begin{split} |J_{2}| &\leq C \frac{1}{R^{2}} \int_{\hat{Q}_{R,t}} |u| \chi_{R}^{q-1} \, d(x,t) \\ &\leq C \frac{1}{R^{2}} \Big( \int_{\hat{Q}_{R,t}} |u|^{p} \chi_{R}^{q} \, d(x,t) \Big)^{1/p} \Big( \int_{\hat{Q}_{R,t}} 1 \, d(x,t) \Big)^{1/q} \\ &\leq C I_{R,t}^{\frac{1}{p}} \frac{1}{R^{2}} R^{\frac{n}{q}} \Big( \int_{\frac{R^{2}}{2}}^{R^{2}} 1 \, dt \Big)^{1/q} \leq C I_{R,t}^{\frac{1}{p}} R^{\frac{n+2}{q}-2} \leq C I_{R,t}^{\frac{1}{p}}. \end{split}$$

Putting the estimates of  $J_1$ ,  $J_2$  and  $J_3$  together we obtain

$$I_R \leq C I_{R,t}^{1/p} + I_{R,x}^{1/p} \rightarrow 0$$

as  $R \to \infty$ . Hence  $u \equiv 0$ , a contradiction with the assumptions on initial data.

# Global existence small data solutions

#### Theorem 2

Consider the Cauchy problem

$$u_{tt} - \bigtriangleup u + u_t = |u|^p, \quad u(0, x) = \varphi(x), \quad u_t(0, x) = \psi(x).$$

Let  $n \leq 4$ ,  $p_{GN}(n) = \frac{n}{n-2}$  for  $n \geq 3$  and let

$$\begin{cases} p > p_{Fuj}(n) & \text{if } n = 1, 2, \\ 2 \le p \le 3 = p_{GN}(3) & \text{if } n = 3, \\ p = 2 = p_{GN}(4) & \text{if } n = 4. \end{cases}$$

Let  $\mathcal{A} := (H^1(\mathbb{R}^n) \cap L^1(\mathbb{R}^n)) \times (L^2(\mathbb{R}^n) \cap L^1(\mathbb{R}^n))$ . Then the following statement holds with a suitable constant  $\varepsilon_0 > 0$ : if  $\|(\varphi, \psi)\|_{\mathcal{A}} \le \varepsilon_0$ , then there exists a unique globally (in time) energy solution  $u \in C([0, \infty), H^1(\mathbb{R}^n)) \cap C^1([0, \infty), L^2(\mathbb{R}^n))$ .

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#### Remark

There exists a constant C > 0 such that the solution and its energy terms satisfy the same decay estimates as the solutions to the associate linear problem

$$\begin{split} \| u(t,\cdot) \|_{L^2} &\leq C(1+t)^{-\frac{n}{4}} \| (\varphi,\psi) \|_{\mathcal{A}}, \\ \| \nabla u(t,\cdot) \|_{L^2} &\leq C(1+t)^{-\frac{n}{4}-\frac{1}{2}} \| (\varphi,\psi) \|_{\mathcal{A}}, \\ \| u_t(t,\cdot) \|_{L^2} &\leq C(1+t)^{-\frac{n}{4}-1} \| (\varphi,\psi) \|_{\mathcal{A}}. \end{split}$$

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#### Remark

Let us consider the corresponding linear Cauchy problem

 $w_{tt}-\bigtriangleup w+w_t=f(t,x), \ w(0,x)=\varphi(x), \ w_t(0,x)=\psi(x).$ 

Using Duhamel's principle w = w(t, x) can be written as

$$w(t,x) = K_0(t,x) *_{(x)} \varphi(x) + K_1(t,x) *_{(x)} \psi(x) + \int_0^t K_1(t-s,x) *_{(x)} f(s,x) \, ds.$$

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# Proof of Theorem 2: 1/4

The space of energy solutions is

$$X(t) = C([0, t], H^1(\mathbb{R}^n)) \cap C^1([0, t], L^2(\mathbb{R}^n))$$

equipped with the norm

$$\begin{aligned} \|u\|_{X(t)} &:= \sup_{0 \le \tau \le t} \left( (1+\tau)^{\frac{n}{4}} \|u(\tau,\cdot)\|_{L^2} + (1+\tau)^{\frac{n+2}{4}} \|\nabla u(\tau,\cdot)\|_{L^2} \right. \\ &+ (1+\tau)^{\frac{n+4}{4}} \|u_{\tau}(\tau,\cdot)\|_{L^2} \right). \end{aligned}$$

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We define the operator

$$N : u \in X(t) \to Nu := K_0(t, x) *_{(x)} \varphi(x) + K_1(t, x) *_{(x)} \psi(x) + \int_0^t K_1(t - s, x) *_{(x)} |u(s, x)|^p ds.$$

Then we show that the following estimates are satisfied:

$$\|Nu\|_{X(t)} \leq C_0 \|(\varphi,\psi)\|_{\mathcal{A}} + C_1(t)\|u\|_{X(t)}^{\rho}, \\ \|Nu - Nv\|_{X(t)} \leq C_2(t)\|u - v\|_{X(t)} (\|u\|_{X(t)}^{\rho-1} + \|v\|_{X(t)}^{\rho-1}).$$

# Proof of Theorem 2: 2/4

#### Proposition

Let u and v be elements of X(t). Then under the assumptions of Theorem 2 the following estimates hold for j + l = 0, 1:

$$\begin{aligned} (1+t)^{l}(1+t)^{\frac{n}{4}+\frac{j}{2}} \|\nabla^{j}\partial_{t}^{l}Nu(t,\cdot)\|_{L^{2}} &\leq C \|(\varphi,\psi)\|_{\mathcal{A}} + C \|u\|_{X(t)}^{p}, \\ (1+t)^{l}(1+t)^{\frac{n}{4}+\frac{j}{2}} \|\nabla^{j}\partial_{t}^{l}(Nu(t,\cdot) - Nv(t,\cdot))\|_{L^{2}} \\ &\leq C \|u-v\|_{X_{0}(t)}(\|u\|_{X(t)}^{p-1} + \|v\|_{X(t)}^{p-1}). \end{aligned}$$

Here the nonnegative constant C is independent of  $t \in [0, \infty)$ .

The derived estimates for solutions to the linear problem implies

$$\begin{aligned} \|\nabla^{j}\partial_{t}^{l}\mathcal{K}_{0}(t,0,x)*_{(x)}\varphi(x)+\nabla^{j}\partial_{t}^{l}\mathcal{K}_{1}(t,0,x)*_{(x)}\psi(x)\|_{L^{2}} \\ &\leq C\,(1+t)^{-l}(1+t)^{-\frac{n}{4}-\frac{j}{2}}\|(\varphi,\psi)\|_{\mathcal{A}}. \end{aligned}$$

# Proof of Theorem 2: 3/4

Using  $K_1(0, 0, x) = 0$  it follows

$$\begin{split} \nabla^{j}\partial_{t}^{l} \int_{0}^{t} \mathcal{K}_{1}(t-s,0,x) *_{(x)} |u(s,x)|^{p} \, ds \\ &= \int_{0}^{t} \nabla^{j}\partial_{t}^{l} \mathcal{K}_{1}(t-s,0,x) *_{(x)} |u(s,x)|^{p} \, ds, \\ \Big\| \int_{0}^{t} \nabla^{j}\partial_{t}^{l} \mathcal{K}_{1}(t-s,0,x) *_{(x)} |u(s,x)|^{p} \, ds \Big\|_{L^{2}} \\ &\leq C \int_{0}^{\frac{t}{2}} (1+t-s)^{-(\frac{n}{4}+\frac{j}{2}+l)} \||u(s,x)|^{p} \|_{L^{2}\cap L^{1}} \, ds \\ &+ C \int_{\frac{t}{2}}^{t} (1+t-s)^{-\frac{j}{2}-l} \||u(s,x)|^{p} \|_{L^{2}} \, ds. \end{split}$$

We use

$$\begin{aligned} \||u(s,x)|^{p}\|_{L^{1}\cap L^{2}} &\leq C \|u(s,\cdot)\|_{L^{p}}^{p} + \|u(s,\cdot)\|_{L^{2p}}^{p}, \\ \||u(s,x)|^{p}\|_{L^{2}} &\leq C \|u(s,\cdot)\|_{L^{2p}}^{p} & \text{ for all } x \neq 0 \text{ for$$

# Proof of Theorem 2: 4/4

Applying Gagliardo-Nirenberg inequality with  $\theta(q) = \frac{n(q-2)}{2q}$ :

$$\|u(s,\cdot)\|_{L^q}^p \leq C \|u(s,\cdot)\|_{L^2}^{p(1- heta(q))} \|
abla u(s,\cdot)\|_{L^2}^{p\, heta(q)} \leq C \|u\|_{X(s)}^p (1+s)^{-rac{(p-1)n}{2}}.$$

We remark that the restriction  $\theta(p) \ge 0$  implies that  $p \ge 2$ , whereas the restriction  $\theta(2p) \le 1$  implies that  $p \le p_{GN}(n)$ if  $n \ge 3$ . Hence, thanks to  $p > p_{Fuj}(n)$  we conclude

$$\begin{split} \left\| \int_{0}^{t} \nabla^{j} \partial_{t}^{l} \mathcal{K}_{1}(t-s,0,x) *_{(x)} |u(s,x)|^{p} ds \right\|_{L^{2}} \\ &\leq C \|u\|_{X(t)}^{p} \int_{0}^{\frac{t}{2}} (1+t-s)^{-(\frac{n}{4}+\frac{j}{2}+l)} (1+s)^{-\frac{(p-1)n}{2}} ds \\ &+ C \|u\|_{X(t)}^{p} \int_{\frac{t}{2}}^{t} (1+t-s)^{-\frac{j}{2}-l} (1+s)^{-\frac{(2p-1)n}{4}} ds \\ &\leq C (1+t)^{-(\frac{n}{4}+\frac{j}{2}+l)} \|u\|_{X(t)}^{p}. \end{split}$$

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# Open problem 1: Klein-Gordon

For the Cauchy problem to semi-linear Klein-Gordon equation

$$u_{tt} - \bigtriangleup u + m^2 u = |u|^{\rho}, \quad u(0, x) = u_0(x), \quad u_t(0, x) = u_1(x),$$

Lindblad-Sogge <sup>8</sup> proved global existence in time for small data energy solution for  $n \le 3$  and  $p > 1 + \frac{2}{n}$  (Fujita). For  $n \le 3$  and  $p \le 1 + \frac{2}{n}$  blow-up results are established. In <sup>9</sup> Keel-Tao conjectured that for sufficiently large dimensions the solution for the semi-linear Cauchy problem has a blow-up for  $p \le 1 + \frac{2}{n} + \epsilon$ ,  $\epsilon > 0$ .

 $^{8}$ H. Lindblad and D. Sogge. Restriction theorems and semilinear Klein-Gordon equations in (1 + 3)-dimensions. Duke Math. J. 85 (1996) 227 – 252.

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Question: What about the critical exponent for large dimensions?

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# Open problem 2: Nakao problem

Let us consider the Cauchy problem for the weakly coupled system of wave equations

$$u_{tt} - \bigtriangleup u + bu_t = |v|^p, \quad v_{tt} - \bigtriangleup v = |u|^q.$$
(2)

If b = 0, the critical behavior (a curve in the p-q plane) of this system is described by the relation

$$\max\left\{\frac{q+2+p^{-1}}{pq-1};\frac{p+2+q^{-1}}{pq-1}\right\} = \frac{n-1}{2}.$$
 (3)

If in (3) the left-hand side is smaller than the right-hand side, then we have GESDS. If in (3) the left-hand side is larger or equal than the right-hand side, then we have, in general, blow-up under some restrictions to the data <sup>10</sup>.

<sup>&</sup>lt;sup>10</sup>D. del Santo, V. Georgiev, E. Mitidieri, Global existence of the solutions and formation of singularities for a class of hyperbolic systems, Geometrical Optics and Related Topics, Eds. F. Colombini and N. Lerner, Progress in Nonlinear Differential Equations and Their Applications, 32 (1997).

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#### What about the critical exponent for $b \neq 0$ in (2)?

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# Open problem 3: Damped wave Scale invariant model

There exists a class of damped wave models for which the critical exponent depends somehow on the Fujita exponent and the Strauss exponent as well. This class is described by the following scale-invariant linear damped wave operators:

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where  $\mu > 0$  is a real parameter. It was recently shown by D'Abbicco'14 that  $p_{Fuj}(n)$  is still the critical exponent when  $\mu \ge \mu_0(n)$ , with  $\mu_0(n)$  sufficiently large.

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It seems to be a challenge to determined the critical exponent in the case  $\mu \in (0, \mu_0(n))$ .

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# Open problem 4: Evolution equations with structural damping

Let us consider the CP

 $u_{tt} + (-\triangle)^{\sigma} u + (-\triangle)^{\theta} u_t = |u|^{\rho}, \qquad (u, u_t)(0, x) = (u_0, u_1)(x),$ 

with p > 1 and  $2\theta \in (0, \sigma]$ . Several papers are devoted to determined the critical exponent to this problem: D'Abbicco & Reissig (2014), D'Abbicco & E. (2014), Duong & Kainane & Reissig (2015), D'Abbicco & E. (2017).

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$$p_c \doteq 1 + \frac{2\sigma}{n-2\theta}, \qquad 2\theta < n.$$

What about the critical exponent for  $\sigma < 2\theta \le 2\sigma$ ? ( $\sigma = 1 = \theta$  is well known as viscoelastic damped wave eq.)

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# Thanks for your attention!